Statistical Thermodynamics

Enthalpy, Free Energy, Entropy, and Heat Capacity of Some Hexafluorides of Octahedral Symmetry *

G. NAGARAJAN and DONALD C. BRINKLEY

Department of Physics and Astronomy, Valdosta State College, Valdosta, Georgia, USA

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A detailed analysis of the molecular structural data and infrared absorption and Raman spectra of the hexafluoride of sulfur, selenium, tellurium, molybdenum, technetium, ruthium, rhodium, tungsten, thenium, osmium, iridium, platinum, uranium, neptunium, and plutonium has been made. These molecules, having the greatest number of symmetry elements of all existing molecules, possess an octahedral symmetry with the symmetry point group O_h . They give rise to six fundamental frequencies of which three are allowed in the Raman spectrum, two are allowed in the infrared absorption spectrum, and one is inactive. The inactive mode in normally determined from the overtones and combinations. On the basis of a rigid rotator and harmonic oscillator model, enthalpy, free energy, entropy, and heat capacity for temperatures from 200 °K to 2000 °K have been computed for these molecules. The results are briefly discussed and compared with available experimental data.

Introduction

From the results of electron diffraction studies, EWING and SUTTON¹ suggested an octahedral symmetry for sulfur and selenium hexafluoride, SEIP and Skolevik² for tellurium hexafluoride, and SEIP and SEIP³ for molybdenum and tungsten hexafluoride. The earlier data for sulfur hexafluoride are from Brockway and Pauling 4. The structure of the hexafluorides of molybdenum, tungsten, and uranium was studied by Braune and Pinnow⁵. The most consistent model was an octahedron in which the X-F distances lay along the rectangular axes, and were in the ratio 1.00 to 1.12 to 1.22 for all three compounds. They did not, however, entirely exclude the regular octahedron in which all the X-F distances are equal. However, tungsten hexafluoride belongs to the point group Oh6'7. BAUER⁸ studied uranium hexafluoride and found

that, while his experimental results were to some extent in harmony with those of Braune and Pin-Now⁵, a more careful interpretation appeared to rule out the structure proposed by them, as well as the totally symmetric one. He proposed, instead, a model without a center of symmetry as the one which best fit his data. But, even then, he himself pointed out two major difficulties with this model. Recently, Seip 9 confirmed an octahedral symmetry for uranium hexafluoride. SCHOMAKER, KIMURA and Weinstock 10 determined internuclear distances and confirmed an octahedral symmetry for two hexafluorides of tungsten, osmium, iridium, uranium, neptunium, and plutonium. A brief survey of two structural data was recently reported by Seip 11 for the hexafluorides of tellurium, molybdenum, tungsten, and uranium. Although the structure for the hexafluorides of technetium, ruthenium, and rhodium is well known from their vibrational assign-

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Reprints request to: Prof. Dr. NAGARAJAN, Valdosta State College, Valdosta, Georgia 31601, USA.

1 V. C. EWING and L. E. SUTTON, Trans. Faraday Soc.

59, 1241 [1963]

² H. M. SEIP and R. STØLEVIK, Acta Chem. Seand. 20, 1535 (1966).

3 H. M. Seip and R. Seip, Acta Chem. Seand. 20, 2698 [1966].
4 L. O. Brockway and L. Pauling, Proc. Natl. Acad.

Sci. 19, 68 [1933]

⁵ H. Braune and P. Pinnow, Z. Physik Chem. **B35**, 239 [1937].

- ⁶ R. V. G. Evans and M. W. Lister, Trans. Faraday Soc. 34, 1358 [1938].
- P. W. Allen and L. A. Sutton, Acta Cryst. 3, 4b
- 8 S. H. BAUER, J. Chem. Phys. 18, 27 [1950]; 18, 994

⁹ H. M. SEIP, Acta Chem. Seand. 19 1955 [1965].

V. SCHOMAKER, M. KIMURA, and B. WEINSTOCK, Preliminary results quoted in P/942 of the second U. N. International Conference on the peaceful uses of Atomic Energy, September 1958.

H. M. Seip, Selected Topics in Structural Chemistry edited by P. Anderson, O. Bastiansen, and S. Fur-BERG, page 25, Universitetsforlaget, Oslo 1967.



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ments, the internuclear distances were not determined experimentally. Recently, the Te-F, Ru-F, and Rh-F distances were evaluated as 1.8512 Å. 1.8775 Å, and 1.8738 Å, respectively, from the stretching force constants of these molecules and from the Badger's empirical relation 12 by NAGARA-JAN 13, 14. From high resolution infrared spectra and other molecular structural studies, several investigators 15-20 reported internuclear distances for the hexafluorides of rhenium, osmium, iridium, platinum, uranium, neptunium, and plutonium.

The Raman and infrared absorption spectra of the group VIA hexafluorides of sulfur, selenium, and tellurium were first completely investigated by GAUNT 21, 22. The normal modes of six oscillations (three valency vibrations and three deformation vibrations) for an octahedral molecule of this type have been given ^{23–25}. One fundamental frequency for selenium hexafluoride and three fundamental frequencies for sulfur and tellurium hexafluoride were assigned to the gaseous state from Raman spectra 26, 27. Recently, Claassen, Goodman, Hol-LOWAY, and SELIG²⁸ studied the Raman spectra of these group VIA hexafluorides in the gaseous state by using powerful laser sources and correctly reassigned the fundamentals. The new values of the fundamentals in cm⁻¹ are given in Table 1. The Raman and infrared absorption spectra of the group VIB hexafluorides of molybdenum and tungsten were studied by many investigators. The Raman spectra of molybdenum hexafluoride have been studied in both liquid 19,29 and gaseous 30 states but with poor resolution and without overtones. The

infrared absorption and Raman spectra of this compound and tungsten hexafluoride were studied by GAUNT²¹. Recently, CLAASSEN and his co-workers 28,31 studied the Raman spectra of these compounds in the gaseous state and accurately reassigned the fundamentals. The recent values of the fundamentals for the hexafluorides of molybdenum and tungsten are given in Table 1.

Recently, Weinstock and Goodman 32 presented a detailed review of electronic, infrared absorption, and Raman spectra of all the then existing hexafluorides. The hexafluorides of rhenium, technetium, and osmium were of particular interest because they offered an almost unique opportunity to study the spectra of molecules with vibronic coupling associated with the dynamic Jahn-Teller effect which was used to account for the unusual widths of certain vibrational bands observed by Weinstock and Claassen 33. Although marked broadening of some of the bands was observed, the vibronic subbands 32 were not identifiable in the Raman spectra of rhenium and osmium hexafluoride 31. The infrared spectrum in the gaseous state and Raman spectrum in the liquid state were studied for technetium hexafluoride by Claassen, Selig, and Malm³⁰ who found that technetium hexafluoride exhibited vibronic coupling to an even greater extent than the hexafluorides of rhenium and osmium. The Raman and infrared absorption spectra of rhenium hexafluoride were studied by GAUNT³⁴. CLAASSEN and his associateds 28 studied the Raman spectra of the group VIIB hexafluorides of technetium and rhenium in the gaseous state with powerful laser sour-

¹² R. M. BADGER, J. Chem. Phys. 2, 128 [1934]; 3, 710

¹³ G. NAGARAJAN, Indian J. Pure Appl. Phys. 1, 232 [1963].

¹⁴ G. NAGARAJAN, Indian J. Pure Appl. Phys. 2, 86 [1964]. ¹⁵ H. H. CLAASSEN, J. Chem. Phys. 30, 968 [1959].

¹⁶ B. Weinstock, H. H. Claassen, and J. G. Malm, J.

Chem. Phys. 32, 181 [1960].
H. C. MATTRAW, N. J. HAWKINS, D. R. CARPENTER, and
W. W. SABOL, J. Chem. Phys. 23, 985 [1955].

H. H. CLAASSEN, B. WEINSTOCK, and J. G. MALM, J. Chem. Phys. 25, 426 [1956].

T. G. BURKE, D. F. SMITH, and A. H. NIELSEN, J. Chem. Phys. 20, 447 [1952].

²⁰ J. G. Malm, B. Weinstock, and H. H. Claassen, J. Chem. Phys. 23, 2192 [1955].

J. GAUNT, Trans. Faraday Soc. 49, 1122 [1953].
 J. GAUNT, Trans. Faraday Soc. 51, 893 [1955].

K. W. F. KOHLRAUSCH, Raman-Spektren, Akademische Verlagsgesellschaft Becker and Erler, Leipzig 1943.

²⁴ G. Herzberg, Molecular Spectra and Molecular Structure II. Infrared and Raman Spectra of Polyatomic Molecules, D. Van Nostrand Company, New York 1960.

²⁵ D. F. HEATH and J. W. LINNETT, Trans. Faraday Soc. 45, 264 [1949]

D. M. Yost, C. C. Steffens, and S. T. Gross, J. Chem. Phys. 2, 311 [1934].

²⁷ C. W. GULLIKSEN, J. R. NIELSEN, and A. T. STAIR, Jr., J. Mol. Spectroscopy 1, 151 [1957].

28 H. H. CLAASSEN, G. L. GOODMAN, J. H. HOLLOWAY, and

H. Selig, J. Chem. Phys. 35, 341 [1970].

K. N. TANNER and A. B. I. DUNCAN, J. Amer. Chem. Soc. 73, 1164 [1951].

³⁰ H. H. CLAASSEN, H. SELIG, and J. G. MALM, J. Chem. Phys. 36, 2888 [1962].

H. H. CLAASSEN and H. SELIG, Israel J. Chem. 7, 499 [1969].

B. Weinstock and G. Goodman, Advan. Chem. Phys. 9, 169 [1965].

B. Weinstock and H. H. Claassen, J. Chem. Phys. 31, 262 [1959].

³⁴ J. Gaunt, Trans. Faraday Soc. 50, 546 [1954].

Molecule	$v_1({ m A}_{1 m g})$	$v_2(\mathrm{E_g})$	$v_3(\mathrm{F_{1u}})$	$v_4(\mathrm{F_{1u}})$	$v_5(\mathrm{F}_{2\mathrm{g}})$	$v_6(\mathrm{F_{2u}})$
Sulfur Hexafluoride	773.5	641.7	939	614	525	347
Selenium Hexafluoride	706.9	658.7	780	437	405	264
Tellurium Hexafluoride	697.1	670.3	752	325	314	197
Molybdenum Hexafluoride	741.5	651.6	741.1	264	318	116
Technetium Hexafluoride	712.9	639	748	265	297	145
Ruthenium Hexafluoride	675	624	735	275	283	186
Rhodium Hexafluoride	634	595	724	283	269	192
Tungsten Hexafluoride	771	677.2	711	258	320	127
Rhenium Hexafluoride	753.7	671	715	257	295	147
Osmium Hexafluoride	730.7	668	720	268	276	205
Iridium Hexafluoride	701.7	645	719	276	267	206
Platinum Hexafluoride	656.4	601	705	273	242	211
Uranium Hexafluoride	667.1	532.5	624	186.2	202	142
Neptunium Hexafluoride	654	535	624	198.6	208	164
Plutonium Hexafluoride	628	523	616	206	211	173

Table 1. Fundamental frequencies in cm⁻¹ for fifteen hexafluorides of octahedral symmetry.

ces, observed the Jahn-Teller effect with over-all frequency differences from 6 to 16 cm⁻¹, and correctly assigned the fundamentals, and many overtones and combinations. The new values of their fundamental frequencies are given in Table 1.

The Raman and infrared absorption spectra of ruthenium and rhodium hexafluoride were first studied in both liquid and gaseous states by Wein-STOCK, CLAASSEN and CHERNICK 35 and all the fundamentals except the inactive mode were assigned. The values of the inactive mode for these two hexafluorides were evaluated from the normal coordinate analysis by Nagarajan³⁶. Then the values of all the fundamentals were reassigned by Weinstock and GOODMAN 32 and they are given in Table 1. The Raman and infrared absorption spectra in both the liquid and gaseous states have been studied for osmium 16, platinum 16, and iridium 17 hexafiuoride. These three hexafluorides were reinvestigated in the vapor along with other hexafluorides, their spectra were analyzed, and their fundamentals given in Table 1, were correctly reassigned by Claasen and Selig³¹. The Raman spectra of uranium hexafluoride were studied in both liquid³⁷ gaseous states 18 and their assignments were in good agreement with each other. The Raman and infrared absorption spectra of uranium hexafluoride were studied in both liquid and gaseous states 19,21,38 and its fundamentals, given in Table 1, were assigned. Claassen and his associates 28 reinvestigated the Raman spectrum of uranium hexafluoride along with other hexafluorides in the gaseous state but their assignments of the fundamentals were similar to those of the previous ones. The Raman and infrared absorption spectra of neptunium and plutonium hexafluoride were studied in both liquid and gaseous states by Malm, Weinstock and CLAASSEN 20 and the fundamentals were assigned. Later, the same studies were undertaken for neptunium hexafluoride by GASNER and FRLEC39 and the fundamentals were reassigned. Recently, Wein-STOCK and GOODMAN 32 reinvestigated these two hexafluorides along with other hexafluorides and correctly reassigned the fundamentals as they are given in Table 1. On the basis of these recent vibrational and structural data, it is aimed here to compute the four thermodynamic quantities, namely, enthalpy, free energy, entropy, and heat capacity on the assumption of a rigid rotator and harmonic oscillator model. The results of the present investigation should be very useful for the evaluation of normal frequencies in other related fluorine-containing compounds, and for the interpretation of the results of experimental thermodynamic quantities, particularly, the entropies and heat capacities at normal pressures.

³⁵ B. Weinstock, H. H. Claassen, and C. L. Chernick,

J. Chem. Phys. 38, 1470 [1963].³⁶ G. NAGARAJAN, Indian J. Pure Appl. Phys. 2, 86 [1964]. 37 J. BIGELEISEN, M. G. MAYER, P. C. STEVENSON, and J. TURKEVICH, J. Chem. Phys. 16, 442 [1948].

³⁸ B. Erlec and H. H. Claassen, J. Chem. Phys. 46 4603

E. L. GASNER and B. FRLEC, J. Chem. Phys. 49, 5135 [1968].

Thermodynamic Functions

One of the best applications of the study of infrared absorption and Raman spectra of polyatomic molecules and other molecular structural determinations is that thermodynamic functions can be statistically computed, namely, enthalpy function $(H_0 - H_0^0)/T$, free enthalpy or Gibbs free energy function $(F_0 - H_0^0)/T$, entropy S^0 , and heat capacity C_n^0 . A rigid rotator and harmonic oscillator model is assumed for each molecule, and all four thermodynamic quantities are computed for a gas in the thermodynamic standard gaseous state of unit fugacity (one atmosphere) for the temperature range from 200°K to 2000°K. The vibrational, translational, and rotional contributions to the total thermodynamic quantities would be computed in the following manner:

The contribution due to molecular vibrations is obtained by summing the appropriate harmonic oscillator function G_{ho} from standard tables of functions 40 over all the normal modes of oscillation of the molecule. For the doubly or triply degenerate vibrations, the term must be doubled or tripled as appropriate. Thus, the vibrational contributions will be given as

$$G_{ extsf{v}} = \sum_{i} G_{ extsf{ho}} u_{i}$$

where d_i is the degeneracy, G_{ho} is the harmonic oscillator function, u_i is the internal thermal energy, and the sum covers all the normal modes. The internal thermal energy, $u_i = h c w_i / T$, is calculated for each normal mode, where h is the Planck's constant, c is the velocity of light, w_i is the ith normal mode in cm^{-1} , and T is the absolute temperature. For each normal mode, the harmonic oscillator contributions to the four thermodynamic quantities, namely, C/R, $(H - H_0)/R T$, $-(F - H_0)/R T$, and S/R, are entered in a tabular form from the standard tables of thermodynamic functions 40 for the corresponding values of the internal thermal energies. After summing up all these values under each column, each total is multiplied by the gas constant R in order to obtain the quantities C, $(H-H_0)/T$, $-(F-H_0)/T$, and S. The value under the column of heat capacity, C, is added with 4R in order to get the value of heat capacity at constant pressure, C_n^0 , for the harmonic oscillator approximation at a pressure of one atmosphere. Similarly, the value under the column of enthalpy function, $(H - H_0)/T$ is added with 4R in order to get the value of enthalpy function, $(H_0 - H_0^0)/T$, for the harmonic oscillator approximation at one atmosphere. The value under column S represents the vibrational contribution to the free energy function.

From the molecular structural data, the rotational and translational contributions to the entropy, heat capacity, and free energy function for one mole of a perfect gas at one atmosphere were obtained from the following relations:

$$\begin{split} S_{\rm tr}^0 + S_{\rm r}^0 &= 2.2868 \, (8 \log T + 3 \log M + \\ &\quad + \log I_{xx} \, I_{yy} \, I_{zz} - 2 \log \sigma) - 7.6965, \\ &\quad - (F_{\rm tr}^0 + F_{\rm r}^0 - H_0^0) / T = S_{\rm rt}^0 + S_{\rm r}^0 - 4 \, R \, , \\ &\quad (C_p^0)_{\rm tr} + (C_p^0)_{\rm r} = (H_{\rm tr}^0 + H_{\rm r}^0 - H_0^0) / T = 4 \, R \, . \end{split}$$

Here S, F, C_p , H, tr, r, T, M, and σ stand for entropy, free energy, heat capacity at constant pressure, enthalpy, translational part, rotational part, temperature in degrees Kelvin, total mass of the molecule, and symmetry number of the point group to which the molecule belongs, respectively, and I_{xx} , I_{yy} , and I_{zz} are the principal moments of inertia in atomic mass units times \mathring{A}^2 along the x-axis, y-axis, and z-axis, respectively. The value of the gas constant, R, is 1.9872 cal/deg mole. These contributions due to translation and rotation are added to the contribution due to vibration in order to obtain the total contribution to the free energy function — $(F_0 - H_0^0)/T$, and the entropy, S^0 .

Results

The hexafluorides of sulfur, selenium, tellurium, molybdenum, technetium, ruthenium, rhodium, tungsten, rhenium, osmium, iridium, platinum, uranium, neptunium, and plutonium have 10 different symmetry elements and 48 different symmetry operations, the highest possible number of symmetry operations ever found in a molecule. These molecules, as described earlier, possess an octahedral symmetry with the symmetry point group Oh. Each molecule possessing several planes of symmetry gives rise, according to the relevant symmetry properties and selection rules 24, to 15 vibrational degrees of freedom which, in turn, constitute only 6 genuine normal (fundamental) modes. The 21 Cartesian displacement vectors generate the representation τ in the following manner:

$$\tau = A_{1g} + E_{g} + F_{1g} + 3 F_{1u} + F_{2g} + F_{2u}$$
.

According to the symmetry operations and selec-

⁴⁰ K. S. PITZER, Quantum Chemistry, Prentice-Hall, Inc., New York 1953.

tions rules (see 24), the rotations and translations belong, respectively, to the F_{1g} and F_{1u} representations. After deleting these, we obtain the following list of genuine normal modes, grouped according to the activities of their fundamentals:

$$egin{aligned} au_{ extbf{v}} &= ext{A}_{1 ext{g}}(R;p) + ext{E}\left(R; ext{d}p
ight) + 2 \, ext{F}_{1 ext{u}}(I;\parallel) \ &+ F_{2 ext{g}}(R; ext{d}p) + ext{F}_{2 ext{u}}(ext{inactive}) \end{aligned}$$

where R, I, p, dp, and \parallel stand for Raman active, infrared active, polarized, depolarized, and parallel, prespectively. The gerade modes are only Raman active, while the ungerade ones are only infrared active. None of the bands observed in the Raman spectrum is observed in the infrared absorption spectrum, thereby indicating that Pauli's exclusion principle is obeyed here, as it must be, since the molecule has a center of symmetry. Further, the occurrence of a genuine normal vibration which is completely inactive as a fundamental must be noted. This is a rare phenomenon in the field of molecular spectroscopy but is occasionally encountered in relatively highly symmetrical molecules. The fundamental frequencies in cm⁻¹ for all the 15 hexafluorides are given in Table 1. On the basis of electron diffraction, microwave, and other structural studies described earlier for all these 15 hexafluorides, the following values of the interatomic distances were adopted for the computations of the principal moments of inertia: S-F = 1.58 Å, Te-F = 1.84 Å, Mo-F = 1.83 Å,Se-F = 1.70 Å,Te-F = 1.8512 Å,Ru-F = 1.8775 Å,Rh-F == 1.8738 Å, W-F = 1.83 Å, Re-F = 1.92 Å, Os-F = 1.831 Å, Ir-F = 1.833 Å, Pt-F = 1.829 Å, U-F = 1.994 Å, Np-F = 1.98 Å, and Pu-F = 1.972 Å.The computed values of the principal moments of inertia $I_{xx} = I_{yy} = I_{zz}$ in units AWUÅ² and gcm² are as follows:

Interatomic	$I_{xx} = I$	$yy = I_{zz}$
distances	$AWA Å^2$	$10^{40}\mathrm{gcm^2}$
S-F ₆	189.7264	315.1602
$Se-F_6$	219.6400	364.8506
$Tl-F_6$	257.3056	427.4181
$Mo-F_6$	254.5164	422.7848
Te-F_{6}	260.4475	432.6372
$Ru-F_6$	267.9005	445.0176
$Rh-F_6$	266.8456	443.2652
$W-F_6$	254.5164	422.7848
Re-F_{6}	280.1664	465.3928
$Os-F_6$	254.7946	423.2469
$Ir-F_6$	255.3516	424.1722
Pt-F6	254.2383	422.3229
$U-F_6$	302.1787	501.9581
$Np-F_6$	297.9504	494.9343
Pu-F6	295,5476	490.9429

Assumed in the computations were a symmetry number of 24, a singlet ground electronic state, and chemical atomic weights 41. Neglected in the calculations were the contributions due to the centrifugal distortion, isotopic mixing, nuclear spins, and interaction between vibration and rotation, since the contributions of these are negligibly small compared to the total thermodynamic quantities due to vibration, rotation, and translation. The computed values of all four thermodynamic quantities in calories per degree mole are given in Table 2 for the 15 hexafluorides.

It is seen from the results for these hexafluorides that the values of heat capacity increase rapidly from 200° to 1000°K, and beyond 2000°K the increase with temperature is negligibly small. Similarly, the values of entropy increase rapidly from 200° to 1500°K, and beyond 1500°K they increase slowly but uniformly. While the values of the enthalpy and free enthalpy functions increase very rapidly near and a little above room temperature, they increase slowly but uniform at higher temperatures. The values of all four thermodynamic quantities at every temperature are, as expected, in increasing order from sulfur to selenium hexafluoride and from selenium to tellurium hexafluoride. This clearly shows that the replacement of a central atom with an atom of higher atomic weight causes lower fundamental frequencies, and correspondingly higher thermodynamic quantities. This is true for the compounds of A-group elements of the periodic table. But, this situation is not observed for many compounds of B-group elements. As an example, all four thermodynamic quantities at a particular temperature for tungsten hexafluoride are not greater than those of molybdenum hexafluoride. Some are greater and some are more or less similar. The reason is that all the observed fundamental frequencies of tungsten hexafluoride are not smaller than those of molybdenum hexafluoride. Actually, some are smaller, some are greater, some are in the same range, and finally all of them are in an irregular order (see Table 1). The same observation may be made for the group VIIB hexafluorides, namely, those of technetium and rhenium, and for the group VIII hexafluorides, namely, those of ruthenium and osmium, and rhodium and iridium. Similarly, the

⁴¹ R. C. Weast, Handbook of Chemistry and Physics, 49th edition, The Chemical Rubber Company, Cleveland, Ohio 1968-1969.

Table 2. Enthalpy, free energy, entropy, and heat capacity of fiften hexafluorides for the ideal gaseous state at a pressure of one atmosphere. (All the quantities are in cal/deg mole.)

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700 22.9611 71.6640 94.6251 33.7129 1700 32.4367 101.4771 139.91 800 24.3842 74.8676 99.2518 34.5713 1800 32.7412 109.5656 142.30 900 25.5306 77.7717 103.3023 35.1689 1900 32.9705 111.2506 144.22 1000 26.5145 80.4810 106.9955 35.6200 2000 33.2286 113.2206 146.44 1100 27.3612 83.0816 110.4428 35.9569 Molybdenum hexafluoride	
800 24.3842 74.8676 99.2518 34.5713 1800 32.1412 109.5656 142.30 900 25.5306 77.7717 103.3023 35.1689 1900 32.9705 111.2506 144.22 1000 26.5145 80.4810 106.9955 35.6200 2000 33.2286 113.2206 146.44 1100 27.3612 83.0816 110.4428 35.9569 Molybdenum hexafluoride	
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1100 27.3612 83.0816 110.4428 35.9569 Molybdenum hexafluoride	
	92 37.4324
1200 28.0721 85.4054 113.4775 36.2207 200 15.6811 57.6675 73.34	86 23.8685
1300 28.7254 87.7866 116.5120 36.4276 273.16 18.4280 62.9703 81.39	
1400 29.2801 89.8659 119.1460 36.6356 298.16 19.2519 64.6169 83.86	
1500 29.7929 92.0221 121.8150 36.8437 300 19.2920 64.6923 83.98	43 28.8266
1600 30.2232 93.9244 124.1476 36.8466 400 22.0892 70.6540 92.74	32 31.8201
1700 30.6386 95.8579 126.4965 36.9439 500 24.8835 75.8835 100.11	41 33.6082
1800 30.9520 97.4321 128.3841 37.0145 600 25.8917 80.4082 106.29	99 34.7228
1900 31.2804 99.1888 130.4692 37.0831 700 27.2209 84.5183 111.73	
2000 31.4368 100.2741 131.7109 37.1118 800 28.2903 88.2446 116.34	
Selection Long House ide 900 29.1744 91.7773 120.95	17 36.3217
Selenium hexafluoride 1000 29.8333 94.5268 124.36	
200 12.1359 53.5166 65.6525 20.0569 1100 30.5034 97.6731 128.17	
273.16 14.9742 57.7284 72.7026 25.0829 1200 31.0123 100.1727 131.18	
298.16 15.8800 59.0843 74.9643 26.4118 1300 31.4902 102.7536 134.24	
300 15.9239 59.1589 75.0828 26.4841 1400 31.8801 105.0491 136.92 400 19.0931 64.2017 83.2948 30.2834 1500 32.2732 107.5590 139.83	92 37.1399 12 37.2215
	33 37.2215 33 37.2782
700 25.1181 76.6435 101.7616 34.9025 1800 33.1114 113.5794 146.69 800 26.3746 80.0754 106.4500 35.5245 1900 33.3129 114.9753 148.28	
900 27.4024 83.2066 110.6090 35.9607 2000 33.5281 117.0143 150.54	
	21 01.1000
1100 20 0429 29 0729 119 0166 26 5297	
1900 90 6567 01 4292 121 0050 26 7215 200 15.4084 57.4430 72.85	
1200 20 2016 02 2022 124 1154 26 2752 275.10 18.2582 02.0990 80.95	
1400 20 6797 06 0400 196 7126 26 0276 298.10 19.0989 04.3193 83.41	
1500 31 1081 98 2457 129 3538 37 0724 300 19.1525 64.4144 83.50	
1600 31 4994 100 2781 131 7775 37 1695 400 21.9954 70.3554 92.35	
1700 31 8316 102 2543 134 0859 37 2331 500 24.1791 75.5067 99.68	
1800 39 1405 104 0705 136 9110 37 9003 000 20.0450 00.0199 105.00	
1900 32.3853 105.6650 138.0503 37.3334 000 27.1575 54.1552 111.55	
1000 20 2526 04 2246 124 12	
Tellurium hexafluoride 1100 29.5320 94.2840 124.13 1100 30.4631 97.1232 127.58	
200 14.2678 56.0095 70.2773 22.6971 1200 31.0018 99.9060 130.90	
273.16 16.9771 60.8465 77.8256 26.9823 1300 31.4741 102.4424 133.91	
298.16 17.8712 62.3674 80.2385 28.1007 1400 31.8901 104.8210 136.71	
300 17.9441 62.4948 80.4389 28.2404 1500 32.2555 107.0994 139.35	
400 20.9390 68.0670 89.0060 31.4197 1600 32.5492 109.0070 141.55	
500 23.2364 72.9803 96.2167 33.3387 1700 32.8616 111.3144 144.17	
600 25.0475 77.4493 102.4968 34.5398 1800 33.0613 112.6910 145.75	
700 26.4402 81.3873 107.8275 35.3067 1900 33.3185 114.8252 148.14	37.4226
800 27.6406 85.1389 112.7795 35.8627 2000 33.5349 116.7560 150.29	09 37.4552

	E.	$-H_0{}^0)/T$				E.	$-H_0^0)/T$		
$T({}^{\circ}{ m K})$	$-H_0{}^0)/T$	100				$H_0^0)/T$	001		
9)	I_0	- P				I_0	- P		
	-				$\widehat{}$	<i>I</i> -	1		
	1	$(F_0$		0	${\bf \mathring{X}}$	1	$(F_0$		0
	$(H_0$	I	S_0	C_p^0	$T({}^{\circ}{ m K})$	$(H_0 -$	Ī	S_0	C_p^0
	Ruth	henium hexofli	uoride		1100	30.4458	98.0157	128.4615	36.7704
200	14.9240	56.4190	71.3430	23.8264	1200	31.0312	101.0268	132.0580	36.9324
273.16	17.8946	61.5330	79.4276	27.8727	1300	31.4608	103.3723	134.8331	37.0447
298.16	18.7771	63.1353	81.9124	28.9264	1400	$31.8505 \\ 32.2156$	105.6323	137.4828	37.1403
300	18.8391	63.2563	82.0954	28.9963	1500 1600	32.2136 32.5472	$107.9324 \\ 110.2466$	$\frac{140.1480}{142.7938}$	$37.2180 \\ 37.2829$
400	21.7871	69.0883	90.8754	31.9966	1700	32.3472 32.8187	111.8919	142.7938 144.7106	37.3382
500	23.9888	74.2315	98.2203	33.7249	1800	33.0466	113.7847	146.8313	37.3773
600	25.7498	78.8031	104.5529	34.8436	1900	33.2835	115.4073	148.6908	37.4179
700 800	27.1212 28.2139	$82.9127 \\ 86.6303$	$\frac{110.0339}{114.8442}$	35.5578	2000	33.4912	117.3472	150.8384	37.4483
900	29.0806	89.8712	118.9518	36.0429 36.3805					
1000	29.8148	92.9033	122.7181	36.6329		7	Ohaniama hana		
1100	30.4555	95.9321	126.3876	36.8177		1	Rhenium hexaj	iuoriae	
1200	30.9897	98.5937	129.5834	36.9626	200	15.4581	58.3207	73.7788	23.9925
1300	31.4646	101.0683	132.5329	37.0792	273.16	18.3151	63.6185	81.9336	27.8360
1400	31.8615	103.4661	135.3276	37.1667	298.16	19.1434	65.2234	84.3668	28.8753
1500	32.2281	105.6635	137.8916	37.2456	300	19.2229	65.3828	84.6057	28.9637
1600	32.5259	107.6330	140.1589	37.3034	400	22.0481	71.2808	93.3289	31.9129
1700	32.8113	109.6308	142.4421	37.3547	500	24.1805	76.4007	100.5812	33.6625
1800 1900	$33.0405 \\ 33.2990$	$111.3765 \\ 113.4290$	$\frac{144.4170}{146.7280}$	$37.3930 \\ 37.4322$	600	25.8754	80.9950	106.8708	34.7772
2000	33.5061	115.4290 115.2168	140.7280 148.7229	37.4522 37.4627	700	27.1957	85.0814	112.2771	35.6131
2000	00.0001	110.2100	140.7220	37.4027	800	28.2795	88.8902	117.1697	35.9863
	Rhc	odium hexaflu	oride		900 1000	29.1777 29.8943	$92.3228 \\ 95.3613$	$\frac{121.5005}{125.2556}$	$36.3468 \\ 36.5975$
					1100	30.4732	98.0896	128.5628	36.7820
200	14.9786	56.4062	71.3848	24.0871	1200	31.0324	100.7564	131.7888	36.9409
273.16 298.16	18.0025	61.5469	79.5494	26.3710	1300	31.5137	103.4845	134.9982	37.0618
300	18.8761 18.9645	$63.1351 \\ 63.2767$	82.0112 82.2412	$\begin{array}{c} 29.1857 \\ 29.2771 \end{array}$	1400	31.9308	105.9006	137.8314	37.1573
400	21.9394	69.1527	91.0921	32.2170	1500	32.2468	107.9225	140.1693	37.2269
500	24.2159	74.3784	98.5943	33.9405	1600	32.5677	110.0554	142.6231	37.2923
600	25.9014	78.8985	104.7999	34.9712	1700	32.8282	112.0632	144.8914	37.3386
700	27.2744	83.0736	110.3480	35.6596	1800 1900	33.0601	113.5325	146.5926	37.3835
800	28.3221	86.6413	114.9634	36.1155	2000	$33.3177 \\ 33.5218$	$\frac{115.6636}{117.1214}$	$\frac{148.9813}{150.6432}$	37.4234 37.4572
900	29.2119	90.0506	119.2625	36.4431	2000	33.3216	117.1214	100.0402	31.4312
1000	29.9237	93.0496	122.9733	36.6798					
1100	30.5700	96.1115	126.6815	36.8603		Osi	nium hexaflue	ride	
1200 1300	$31.1077 \\ 31.5707$	98.8008 101.3474	$\frac{129.9085}{132.9181}$	37.0013 37.1110			,		
1400	31.9564	103.5305	135.4869	37.1110	200	14.7831	57.1297	71.9128	23.6509
1500	32.3308	105.9775	138.3083	37.2712	273.16	17.7368	62.1956	79.9324	27.6712
1600	32.6418	108.1334	140.7752	37.3261	298.16	18.6185	63.7864	82.4049	28.7382
1700	32.9093	109.9722	142.8815	37.3757	$\frac{300}{400}$	$18.6780 \\ 21.6212$	$63.9012 \\ 69.6956$	$82.5792 \\ 91.3168$	28.8055 31.8336
1800	33.1505	111.9261	145.0766	37.4117	500	21.0212 23.8671	74.7889	98.6560	33.6386
1900	33.3815	113.6332	147.0147	37.4489	600	25.5932	79.3096	104.9028	34.7493
2000	33.5893	115.4232	149.0125	37.4795	700	26.9493	83.3227	110.2720	35.4797
	<i>m</i>				800	28.0512	86.9743	115.0255	35.9823
	Tui	ngsten hexaflu	oriae		900	28.9323	90.2916	119.2239	36.3241
200	15.5523	58.4007	73.9530	23.8520	1000	29.6871	93.4221	123.1092	36.5840
273.16	18.3070	63.6450	81.9520	27.7118	1100	30.3403	96.2939	126.6342	36.7857
298.16	19.1706	65.3375	84.5081	28.7626	1200	30.8790	98.9175	129.7965	36.9356
300	19.2072	65.4091	84.6163	28.8212	1300	31.3106	101.2240	132.5346	37.0455
400 500	$22.0030 \\ 24.4623$	$71.3058 \\ 77.0764$	93.3088 101.5387	$31.8110 \\ 33.7143$	1400 1500	$31.7877 \\ 32.1476$	$104.1915 \\ 106.1503$	$\begin{array}{c} 135.9792 \\ 138.2979 \end{array}$	$37.1478 \\ 37.2267$
600	24.4623 25.8589	81.1255	101.5387 106.9844	33.7143 34.7389	1600	$32.1476 \\ 32.4455$	106.1503	138.2979 140.6144	37.2267 37.2847
700	27.1752	85.1643	112.3395	35.4670	1700	32.4455 32.7295	108.1089 110.5382	140.0144 142.8677	37.3378
800	28.2293	88.8100	117.0393	35.9599	1800	33.0097	112.1809	145.1906	37.3843
900	29.1364	92.3476	121.4840	36.3205	1900	33.1951	113.5522	146.7473	37.4175
1000	29.8760	95.4555	125.3315	36.5874	2000	33.4184	115.3477	148.7661	37.4512

	-	$-H_0^0)/T$				-	$-H_0^0)/T$		
	$-H_0^0)/T$	6				$H_0^0)/T$	6		
	I_0	Н-				I_0	Н -		
Œ	- T				\overline{a}	<i>I</i> –	1		
$T({}^{\circ}{ m K})$	ွ်	$(F_0$	_	0	$T({}^{\circ}{ m K})$	٥	$(F_0$		0
T	$(H_0$	l	S_0	C_p^0	T	$(H_0$	I,	80	C_p^{α}
	Irr	idium hexaflu	oride		1000 1100	31.1015 31.6923	$100.9528 \\ 103.8955$	132.0543 135.5878	36.9338 37.0719
200	14.8095	57.1554	71.9649	23.7634	1200	32.1667	106.7815	138.9482	37.1798
273.16	17.7965	62.2446	80.0411	27.8169	1300	32.5283	109.1355	141.6638	37.2627
298.16	18.6845	63.8397	82.5242	28.8905	1400	32.8596	111.4853	144.3449	37.3278
$\frac{300}{400}$	18.7424 21.7067	$63.9426 \\ 69.7728$	$82.6850 \\ 91.4795$	28.9514 31.9574	1500 1600	$33.1702 \\ 33.4327$	$\frac{113.8590}{115.9720}$	$\frac{147.0292}{149.4047}$	37.3831 37.4278
500	23.9510	74.8806	98.8316	33.7280	1700	33.6734	118.0953	151.7687	37.4648
600	25.6916	79.4315	105.1231	34.8331	1800	33.8918	120.2203	154.1121	37.4946
700	27.0274	83.4400	110.4674	35.5336	1900	34.0907	121.9644	156.0551	37.5236
800	28.1303	87.1086	115.2389	36.0274	2000	34.2580	123.7866	158.0446	37.5444
900	29.0229	90.4379	117.4736	36.3690					
1000 1100	29.7751 30.4110	93.5539 96.4592	$\frac{123.3290}{126.8702}$	$36.6236 \\ 36.8132$		Nept	unium hexafl	uoride	
1200	30.9517	99.0305	129.9822	36.9608	200	17.0369	60.0705	77.1074	26.3482
1300	31.3784	101.4112	132.7896	37.0663	273.16	20.0446	65.8507	85.8953	20.3482 29.9206
1400	31.8505	104.1674	136.0179	37.1679	298.16	20.9295	67.6674	88.5969	30.8305
1500	32.1943	106.2251	138.4194	37.2398	300	20.9736	67.7635	88.7371	30.8587
1600	32.5246	108.4261	140.9507	37.3050	400	23.8113	74.2485	86.0598	33.3612
1700	32.7859	110.0215	142.8074	$37.3507 \\ 37.3964$	500	25.8573	79.6839	105.6412	34.7369
1800 1900	$33.0522 \\ 33.2688$	$\frac{112.2268}{113.8387}$	$\frac{145.2790}{147.1075}$	37.3904 37.4320	600	27.4071	84.6182	112.0253	35.5727
2000	33.4687	115.5644	149.0331	37.4618	700 800	28.6110 29.6052	88.8637 92.9449	117.4747	36.1214
2000	00.2001	110,0011	110,0001	3772320	900	30.3567	92.9449 96.3794	$122.5501 \\ 126.7361$	$36.4720 \\ 36.7345$
	Pla	tinum hexaflu	orida		1000	30.9837	99.4116	130.3953	36.9217
	1 111	iinum nexujiu	orue		1100	31.5624	102.7038	134.2662	37.0633
200	15.1179	57.2923	72.4102	24.3157	1200	32.0051	105.2053	137.2104	37.1722
273.16	18.1510	62.4767	80.6277	28.3179	1300	32.4320	107.9917	140.4237	37.2595
298.16	19.0305	64.0757	$83.1062 \\ 83.3278$	$\frac{29.3431}{29.4238}$	1400	32.7903	110.5107	143.3010	37.3274
$\frac{300}{400}$	$\frac{19.1102}{22.0813}$	64.2176 70.1386	92.2199	32.3163	1500 1600	$33.0655 \\ 33.3280$	$\frac{112.5227}{114.5452}$	$\frac{145.5882}{147.8732}$	$37.3785 \\ 37.4232$
500	23.9709	74.5244	98.4953	33.8896	1700	33.5683	116.5661	150.1344	37.4606
600	26.0241	79.9042	105.9283	35.0354	1800	33.7860	118.5726	152.3586	37.4910
700	27.3553	84.0341	111.3894	35.6952	1900	33.9967	120.6370	154.6337	37.5182
800	28.4158	87.6920	116.1078	36.1482	2000	34.1630	121.7946	155.9576	37.5401
900	29.2525	91.0181	120.2706	36.4660					
1000 1100	$30.0656 \\ 30.6200$	94.3737 96.9863	$\frac{124.4393}{127.6063}$	$36.7154 \\ 36.8793$		Plut	onium hexaflu	ioride	
1200	31.1653	99.8531	131.0184	37.0137	200	16.8660	50 9450	76.7119	26.3911
1300	31.6290	102.4147	134.0437	37.1226	273.16	19.9478	$59.8459 \\ 65.5910$	85.5388	30.0171
1400	32.0232	104.6546	136.6778	37.2118	298.16	20.8305	67.3793	88.2098	30.9173
1500	32.3749	107.0555	139.4304	37.2758	300	20.8760	67.4760	88.3506	30.9708
1600	32.6716	108.4144	141.5850	37.3360	400	23.7360	73.9091	97.6451	33.4170
1700 1800	32.9544 33.2167	$\frac{111.0994}{112.9809}$	$\frac{144.0538}{146.1976}$	$37.3799 \\ 37.4250$	500	25.8251	79.4153	105.2404	34.8073
1900	33.4246	112.3803 114.7238	148.1484	37.4250 37.4562	600	27.4048	84.3379	111.7427	35.6290
2000	33.6338	116.5437	150.1775	37.4840	700 800	$28.6116 \\ 29.5774$	88.6087 92.4972	$\begin{array}{c} 117.2203 \\ 122.0746 \end{array}$	36.1556 36.5127
					900	30.3520	95.9091	126.2611	36.7661
	IIro	anium hexaflu	oride		1000	31.0085	99.2220	130.2305	36.9449
200				20.0500	1100	31.5342	102.0566	133.5908	37.0802
200	$\frac{17.5765}{20.5018}$	60.8668	78.4433 87.2056	26.6560	1200	32.0111	106.9138	136.9249	37.1858
$273.16 \\ 298.16$	20.5018 21.3355	$66.7938 \\ 68.6003$	$87.2956 \\ 89.9358$	30.0956 30.9620	1300	32.4308	107.6472	140.0780	37.2700
300	21.4020	68.7644	90.1664	30.9801	1400	32.7626	109.8869	142.6496	37.3348
400	24.1273	75.2887	99.4160	34.2159	1500 1600	$33.0736 \\ 33.3368$	$\frac{112.1531}{114.1532}$	$\frac{145.2267}{147.4900}$	$37.3898 \\ 37.4344$
500	26.1256	80.8699	106.9955	34.7797	1700	33.5981	116.4313	150.0294	37.4344
600	27.6484	85.8222	113.4706	35.6055	1800	33.8157	118.4227	152.2384	37.5005
700	28.8433	90.2032	119.0465	36.1474	1900	33.9731	120.0098	153.9829	37.5218
800 900	$29.7795 \\ 30.5169$	94.1135 97.5619	$\frac{123.8930}{128.0788}$	$36.5032 \\ 36.7478$	2000	34.1950	122.1544	156.3497	37.5478
900	90.9109	91.5019	120.0100	50.1410					

values of thermodynamic quantities for uranium, neptunium, and plutonium hexafluoride show the same irregular order at every temperature which is due to the observed fundamental frequencies.

Experimental and calculated values for hexafluoride molecules in the gaseous state are compared in Table 3. Encluding the entropies of molybdenum and tungsten hexafluoride the agreement is ex-

Table 3. Comparison of Experimental and Calculated Values (cal/deg mole).

Hexa- fluoride	Quantity	Experi- mental 298 °K	Ref.	Calculated (this work) 298.16 °K
Sulfur	S0	69.5	41	69.7081
Tellurium	S^0	80.67	41	80.2385
Molybdenum	S^0	80.6*	42	83.8688
		80.05	19	
	C_{p}^{0}	28.35	19	28.7561
Tungsten	${}^{C_{m p}}_{S^0}$	81.75	19	84.5081
0	C_{p}^{0} S^{0}	28.44	19	28.7626
Uranium	So	90.76	41	89.9358

^{* 298.15 °}K.

cellent. The small disagreement in the entropy values for these molecules may be attributed to some error in the experimental observation of the molecular structural data. The good agreement between the experimental and spectroscopic values of thermodynamic quantities for some of the molecules studied here shows that the results presented here are very reliable and would be very useful in the future for the interpretation of the results of other experimental thermodynamic quantities and for the evaluation of normal frequencies in other related molecular systems having similar chemical bonds.

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⁴² A. P. Brady, O. E. Myers, and J. K. Clauss, J. Phys. Chem. **64**, 588 [1960].